

Appendix D

Does a normal mode carry Physical Momentum?

Back to our 1D monatomic chain example

Physical momentum = Mass · Velocity

$$= \sum_{\text{all atoms } n} M \cdot \frac{dx_n}{dt} \quad \begin{matrix} (N \text{ atoms}) \\ (N \gg 1) \end{matrix}$$

$$u_n = A e^{iqna} e^{-i\omega t} \equiv A(t) e^{iqna}$$

$$\therefore \text{Physical momentum} = M \frac{dA}{dt} \sum_{n=0}^{N-1} e^{iqna}$$

This is the same as: $\frac{1 - e^{iqNa}}{1 - e^{iqa}}$ (in 1D)

• but now \vec{q} is in 1st B.Z.

Vanishes except $\vec{q} = \vec{G}$

(i) But for any $\vec{q} \neq 0$ in 1st B.Z., $\vec{q} \neq \vec{G}$

\therefore Physical momentum = 0 for any $\vec{q} \neq 0$

The physical reason is that only relative coordinates of the atoms are involved.

(ii) The Only \vec{G} in the 1st B.Z. is $\vec{G} = 0$.

\therefore The $\vec{q} = 0$ mode is an exceptional case.

For $q=0$ (1D chain),

$$\text{physical momentum} = M \frac{dA}{dt} \sum_{n=0}^{N-1} 1$$

$$= MN \frac{dA}{dt}$$

representing a uniform translation of the crystal as a whole.

Remarks:

- Since the physical momentum is different from $\hbar q$ and $\hbar q$ appears in "selection rules" or "momentum conservation" rules, we call $\hbar q$ the crystal momentum of a phonon in the mode \vec{q} (or (s, \vec{q}) where s labels the branches)

- Recall that when the media is dispersive (i.e. $\omega(\vec{q})$), the velocity of a wave packet is the group velocity:

$$\vec{V}_g = \vec{\nabla}_{\vec{q}} \omega(\vec{q}) \leftarrow \text{dispersion relation gradient w.r.t. } \vec{q}$$

$$\frac{\partial \omega}{\partial q} \text{ in 1D}$$